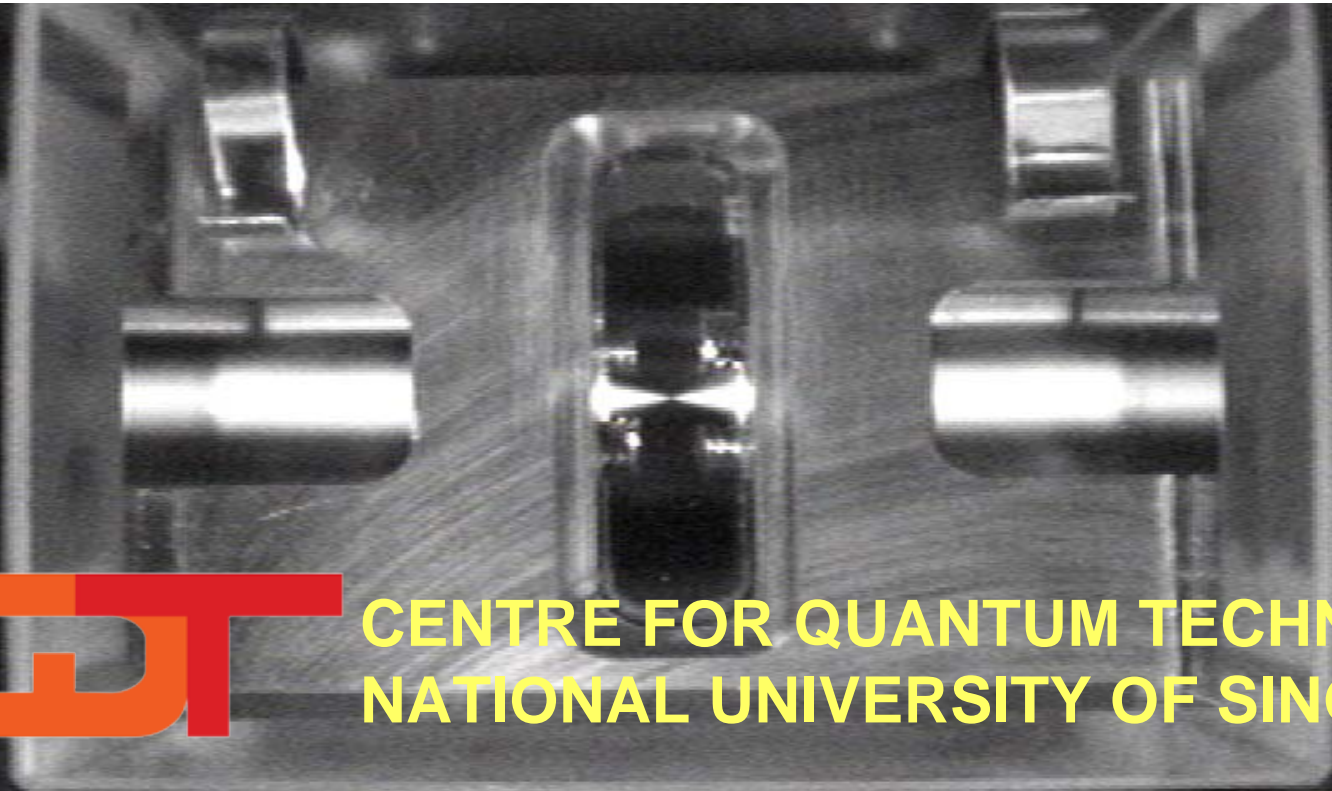


STRONG INTERACTION BETWEEN LIGHT AND SINGLE TRAPPED ATOM WITHOUT A CAVITY



CENTRE FOR QUANTUM TECHNOLOGIES
NATIONAL UNIVERSITY OF SINGAPORE

Meng Khoon Tey, Zilong Chen, Syed Abdullah Aljunid,
Brenda Chng, Gleb Maslennikov and Christian Kurtsiefer

Motivation

Light-Atom Interface for Quantum Information

quantum repeaters & quantum memory

(J. I. Cirac et. al, *PRL*, 78, 3221, L. M. Duan et. al, *Nature*, 414, 413)

conditional phase gates on photonic qubits

(S. M. Savage et. al, *Opt.Lett*, 15, 628)

Precision Spectroscopy with Cold Atoms

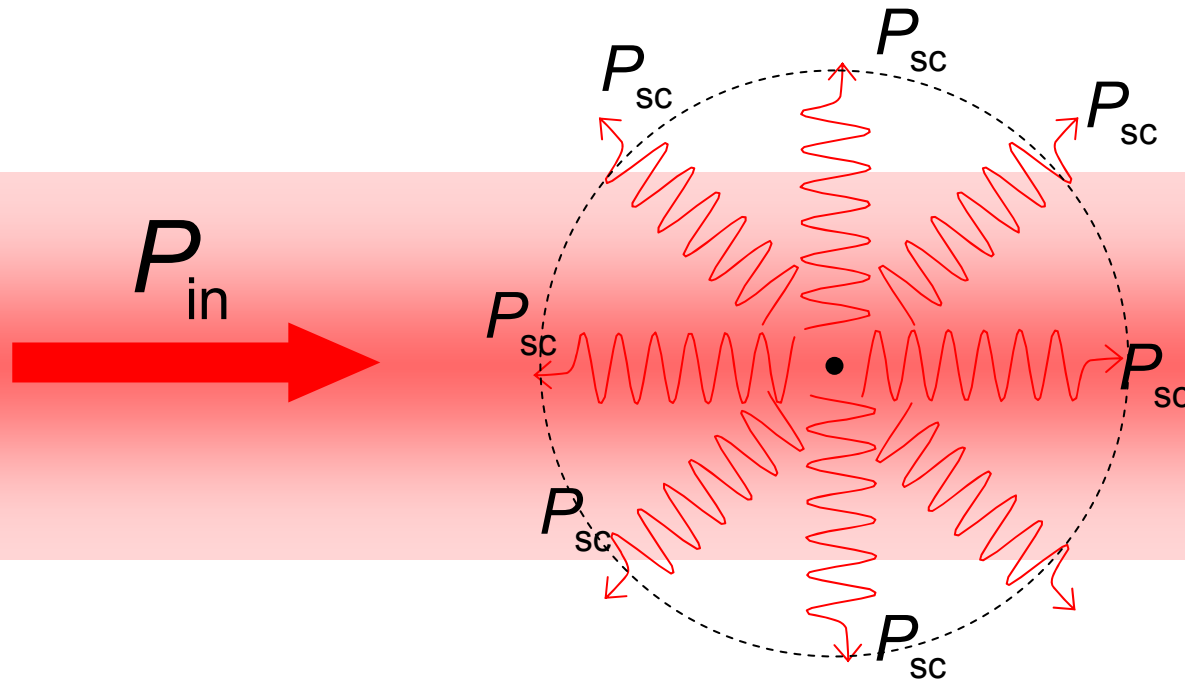
determination of atomic polarizabilities

(B. Arora et. al, *PRA*, 76, 052616, M. Safronova, et. al, *PRA*, 73, 022505)

atomic response to the excitation profile

(P.V. Elyutin, arXiv:quant-ph/0802.0913, E. S. Kyoseva et al, *PRA*, 73, 023420)

Physical System: Atom in a Light Field

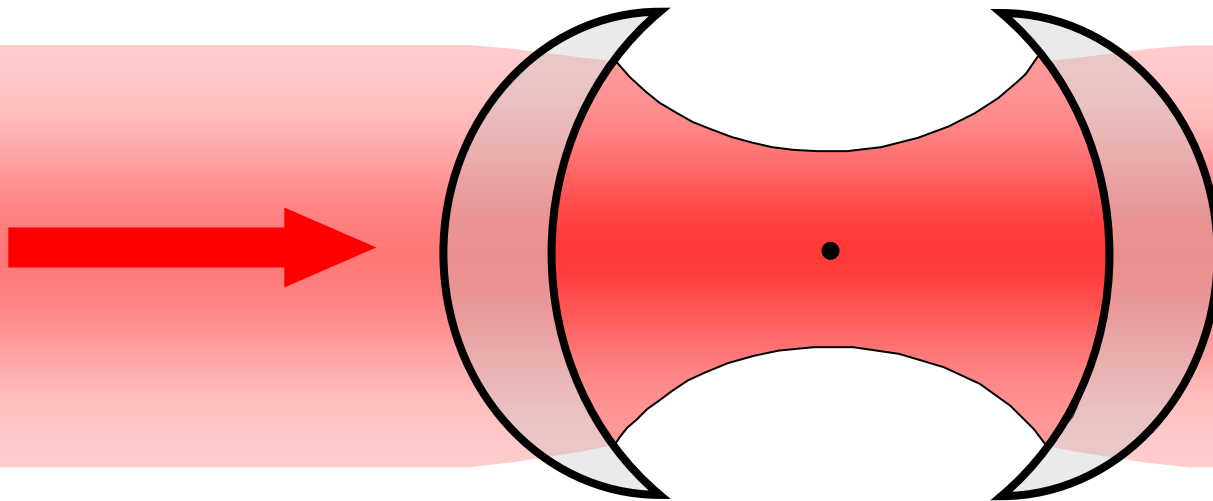


The scattering probability $\rho_{sc} = \frac{P_{sc}^{tot}}{P_{in}}$

Concentration of the incoming field
at the position of the atom is necessary!

Physical System: Atom in a Light Field

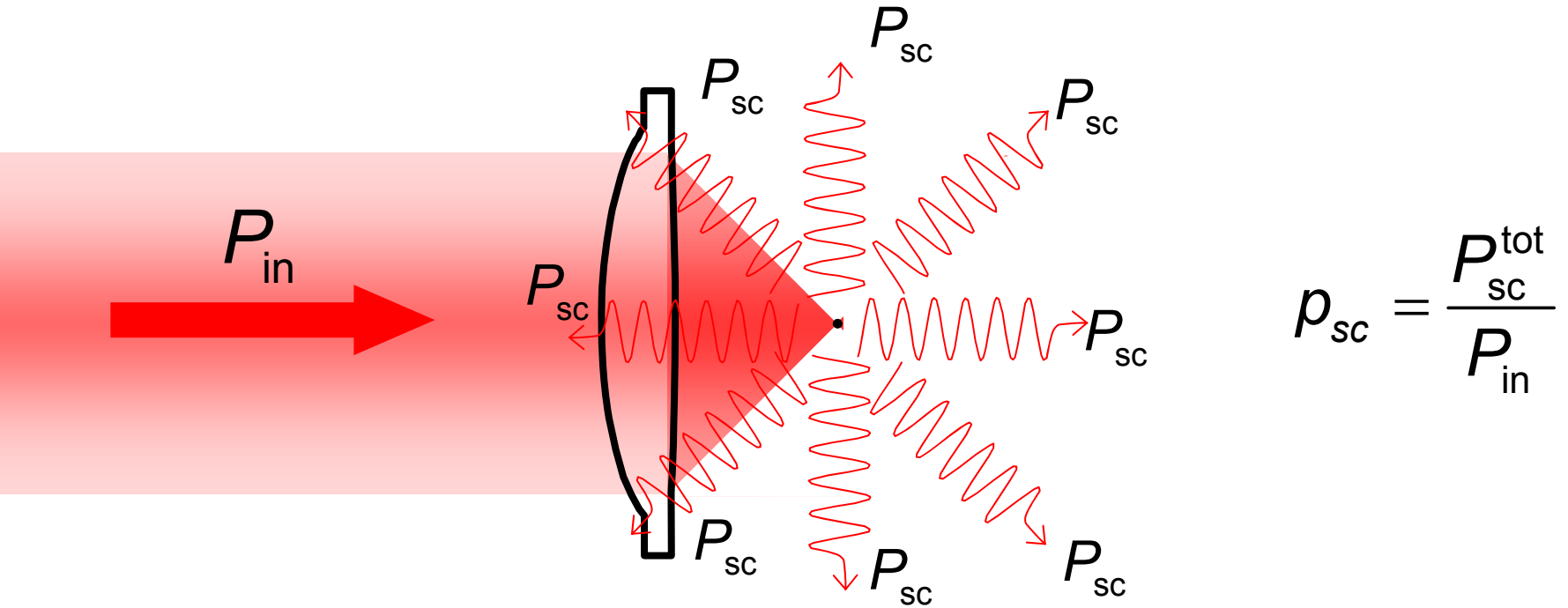
One solution: use a high-finesse cavity around the atom



Many ongoing experiments
CalTech, Univ. of Georgia, Max-Planck-Institute, etc...

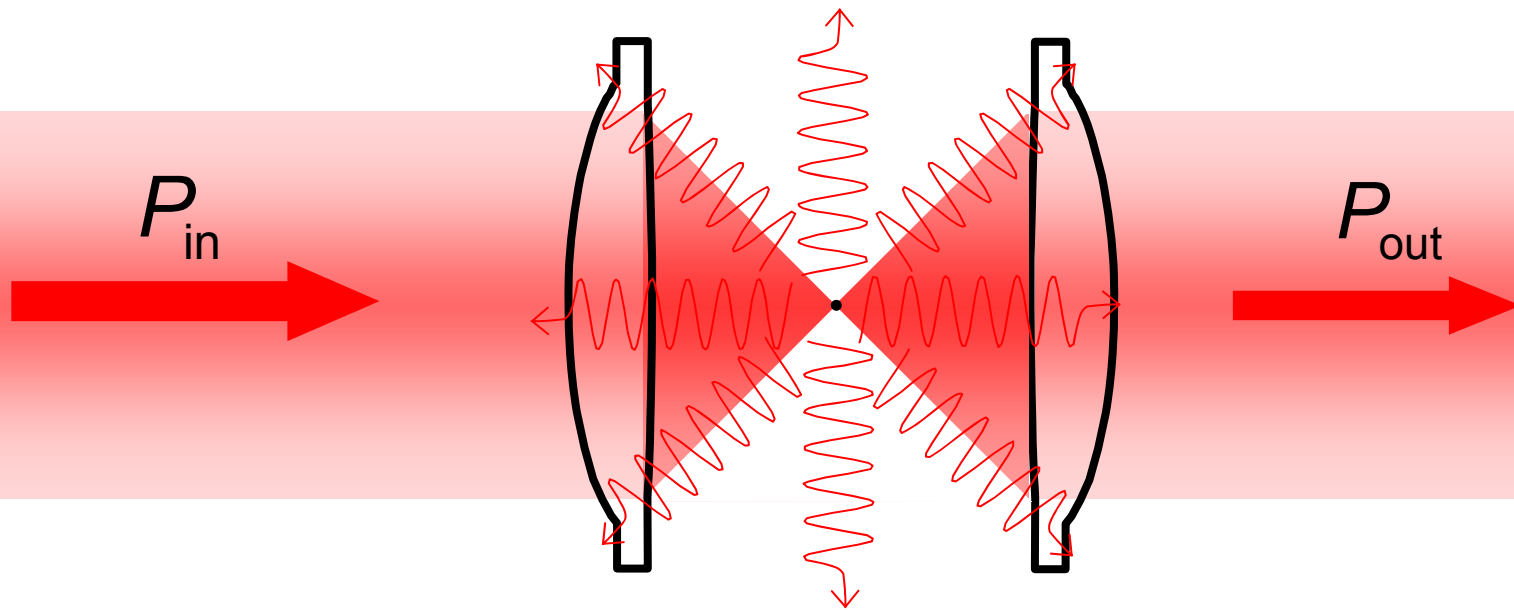
Physical System: Atom in a Light Field

Or just use a lens to focus the beam to the atom



Hard to obtain P_{sc}^{tot} , need a detector, covering 4π solid angle

Physical System: Atom in a Light Field



Relate P_{sc} to the transmission of the probe through the atom

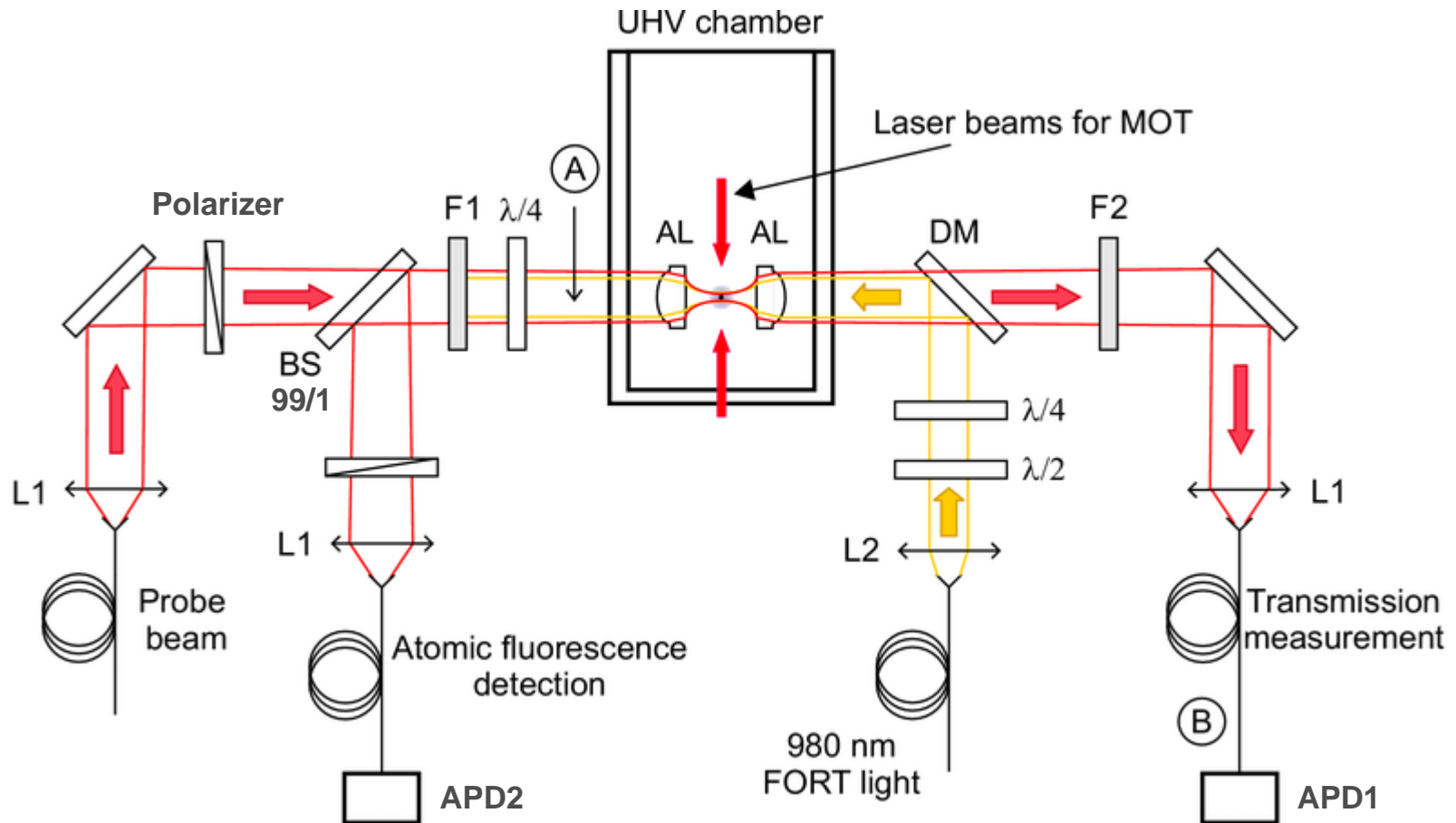
$$T = \frac{P_{out}}{P_{in}} = \frac{P_{in} - P_{sc} + \alpha P_{sc}}{P_{in}}$$

$$P_{sc} = \frac{1 - T}{1 - \alpha}$$

If α (collection efficiency) is small

$$P_{sc} \approx 1 - T$$

Experimental Setup

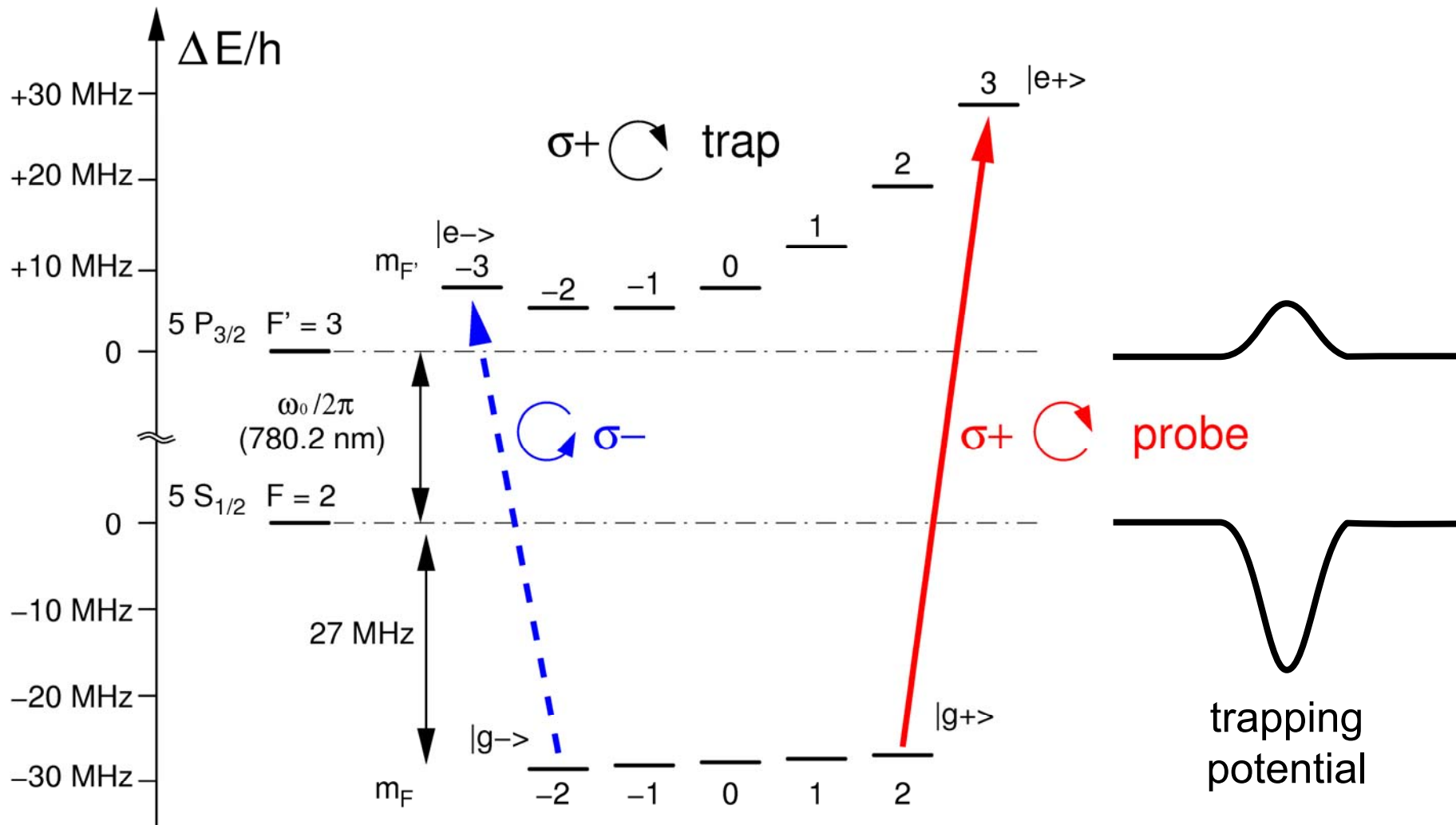


AL – aspheric lenses ($f = 4.5$ mm, full NA = 0.55),
DM – dichroic mirror, F1 – filters to block 980 nm light,

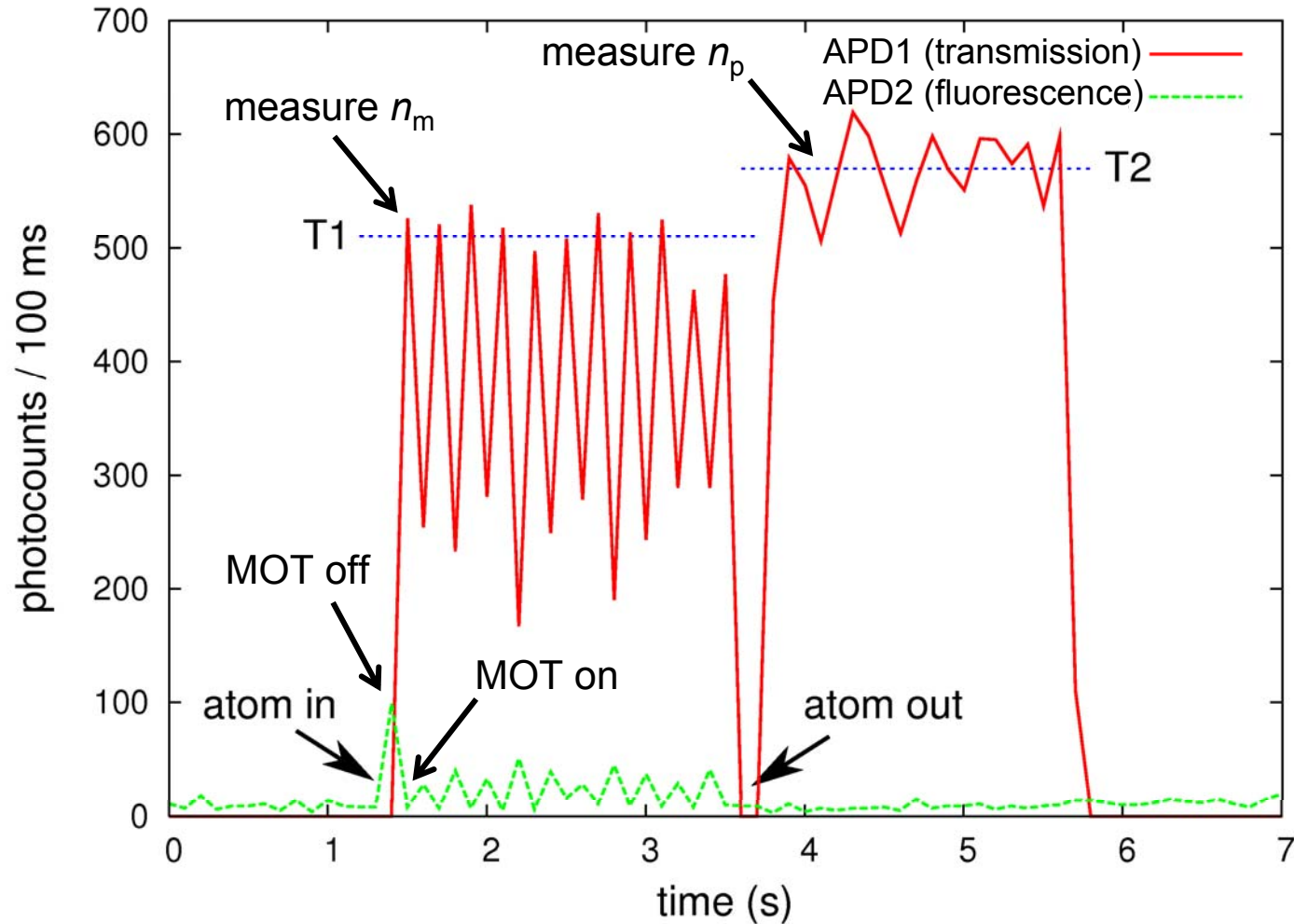
Losses: 55% transmission from (A) to (B)
dominated by reflection loss at optical surfaces

^{87}Rb Atom in the Optical Dipole Trap: AC Stark shift

Dipole trap laser wavelength: 980 nm

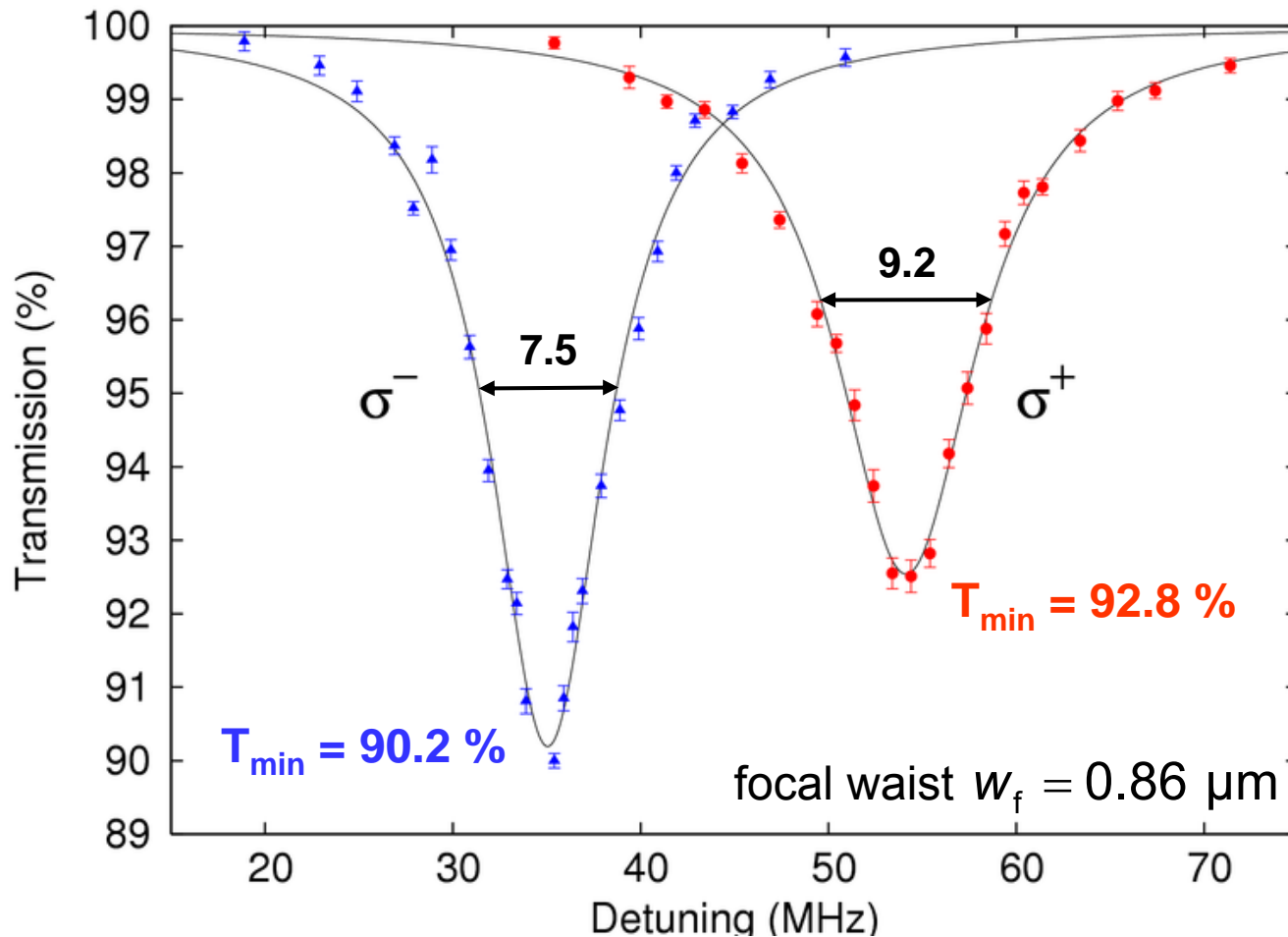


Experimental Procedure



$$T = \frac{T_1}{T_2} = \frac{\sum n_m t_p}{\sum t_m n_p}$$

Experimental Results



Actual scattering rate
2500 photons s^{-1}

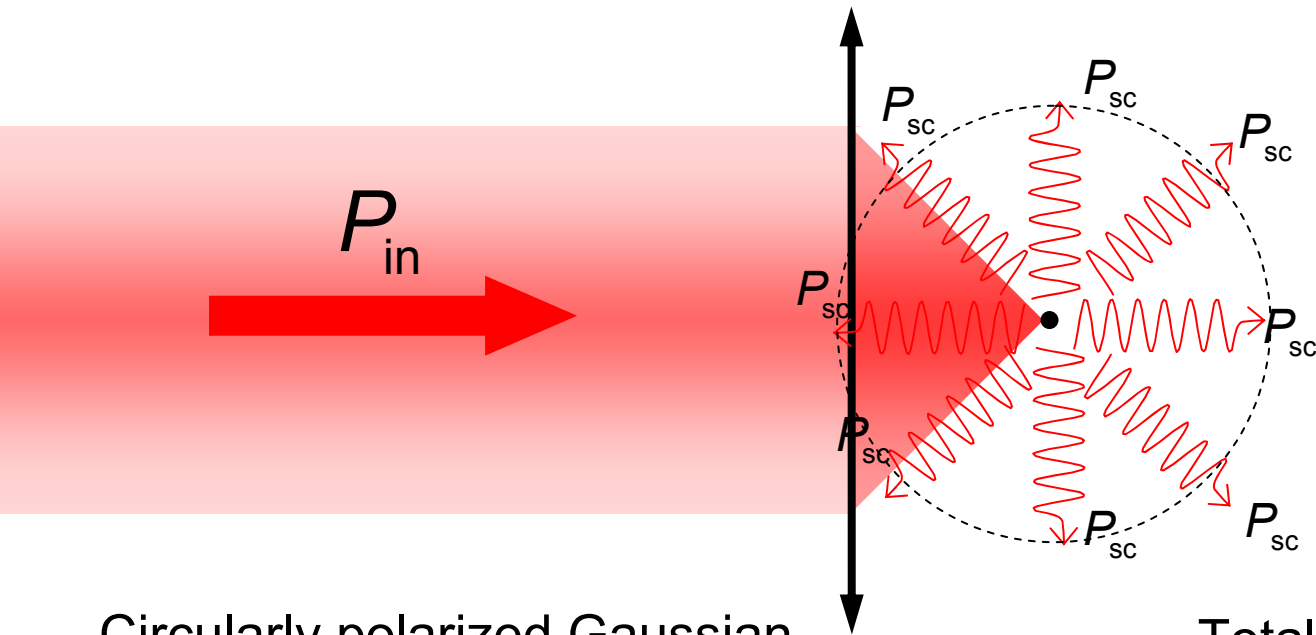
Natural linewidth
6 MHz

Collection efficiency
 $\alpha \sim 5\%$, so $p_{\text{sc}} \approx 9.8\%$

Resonance frequency depends on the AC Stark shift of the energy level

Infer polarizability!

Theoretical model



Circularly polarized Gaussian beam at the lens

$$P_{in} = \frac{1}{4} \epsilon_0 \pi c E_L^2 W_L^2$$

Total power scattered by a two-level atom

$$P_{sc} = \frac{3\epsilon_0 c \lambda^2 E_A^2}{4\pi}$$

$$p_{sc} \equiv \frac{P_{sc}}{P_{in}} = \frac{3\lambda^2}{\pi^2 W_L^2} \left(\frac{E_A}{E_L} \right)^2$$

Theoretical model

$$\rho_{\text{sc}} = \frac{3\lambda^2}{\pi^2 w_L^2} \left(\frac{E_A}{E_L} \right)^2$$

1. Paraxial approximation

$$\left(\frac{E_A}{E_L} \right)^2 = \left(\frac{w_f}{w_L} \right)^2 \quad \rho_{\text{sc}} = \frac{3\lambda^2}{\pi^2 w_f^2}$$

Theoretical model

$$\rho_{\text{sc}} = \frac{3\lambda^2}{\pi^2 W_L^2} \left(\frac{E_A}{E_L} \right)^2$$

1. Paraxial approximation $\left(\frac{E_A}{E_L} \right)^2 = \left(\frac{W_f}{W_L} \right)^2$ $\rho_{\text{sc}} = \frac{3\lambda^2}{\pi^2 W_f^2}$

Inappropriate for strong focusing regime

Theoretical model

$$\rho_{\text{sc}} = \frac{3\lambda^2}{\pi^2 w_L^2} \left(\frac{E_A}{E_L} \right)^2$$

1. Paraxial approximation $\left(\frac{E_A}{E_L} \right)^2 = \left(\frac{w_f}{w_L} \right)^2$ $\rho_{\text{sc}} = \frac{3\lambda^2}{\pi^2 w_f^2}$

2. Decomposition of the field into modes with cylindrical symmetry and model the action of the lens as a local phase shifter of the wavefront curvature (S. J. van Enk and H. J. Kimble, PRA, **63**, 023809)

Theoretical model

$$\rho_{\text{sc}} = \frac{3\lambda^2}{\pi^2 W_L^2} \left(\frac{E_A}{E_L} \right)^2$$

1. Paraxial approximation $\left(\frac{E_A}{E_L} \right)^2 = \left(\frac{W_f}{W_L} \right)^2$ $\rho_{\text{sc}} = \frac{3\lambda^2}{\pi^2 W_f^2}$

2. Decomposition of the field into modes with cylindrical symmetry and model the action of the lens as a local phase shifter of the wavefront curvature (S. J. van Enk and H. J. Kimble, PRA, **63**, 023809)

Original paper

parabolic wave front after the lens



**Allows analytical integration
for field decomposition**

**Does not efficiently concentrate
the incoming energy at the focus
for strong focusing regime**

Theoretical model

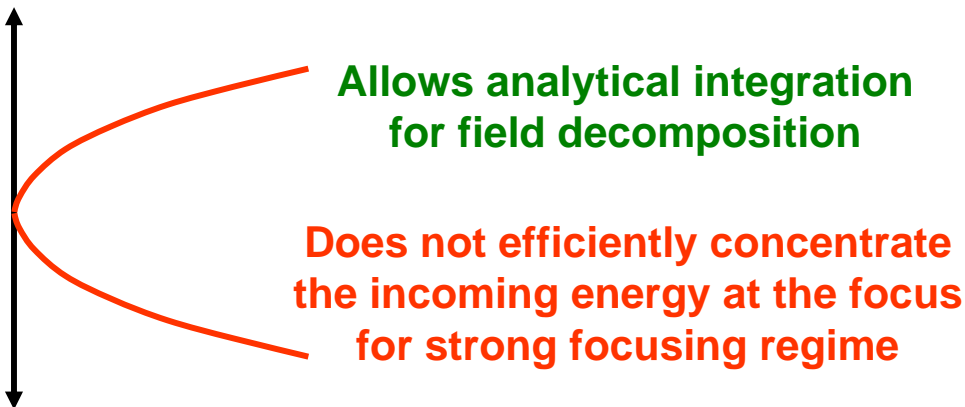
$$\rho_{\text{sc}} = \frac{3\lambda^2}{\pi^2 W_L^2} \left(\frac{E_A}{E_L} \right)^2$$

1. Paraxial approximation $\left(\frac{E_A}{E_L} \right)^2 = \left(\frac{W_f}{W_L} \right)^2$ $\rho_{\text{sc}} = \frac{3\lambda^2}{\pi^2 W_f^2}$

2. Decomposition of the field into modes with cylindrical symmetry and model the action of the lens as a local phase shifter of the wavefront curvature (S. J. van Enk and H. J. Kimble, PRA, **63**, 023809)

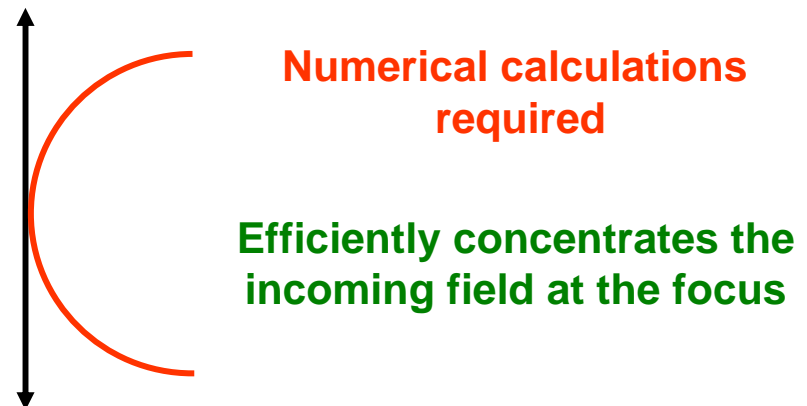
Original paper

parabolic wave front after the lens



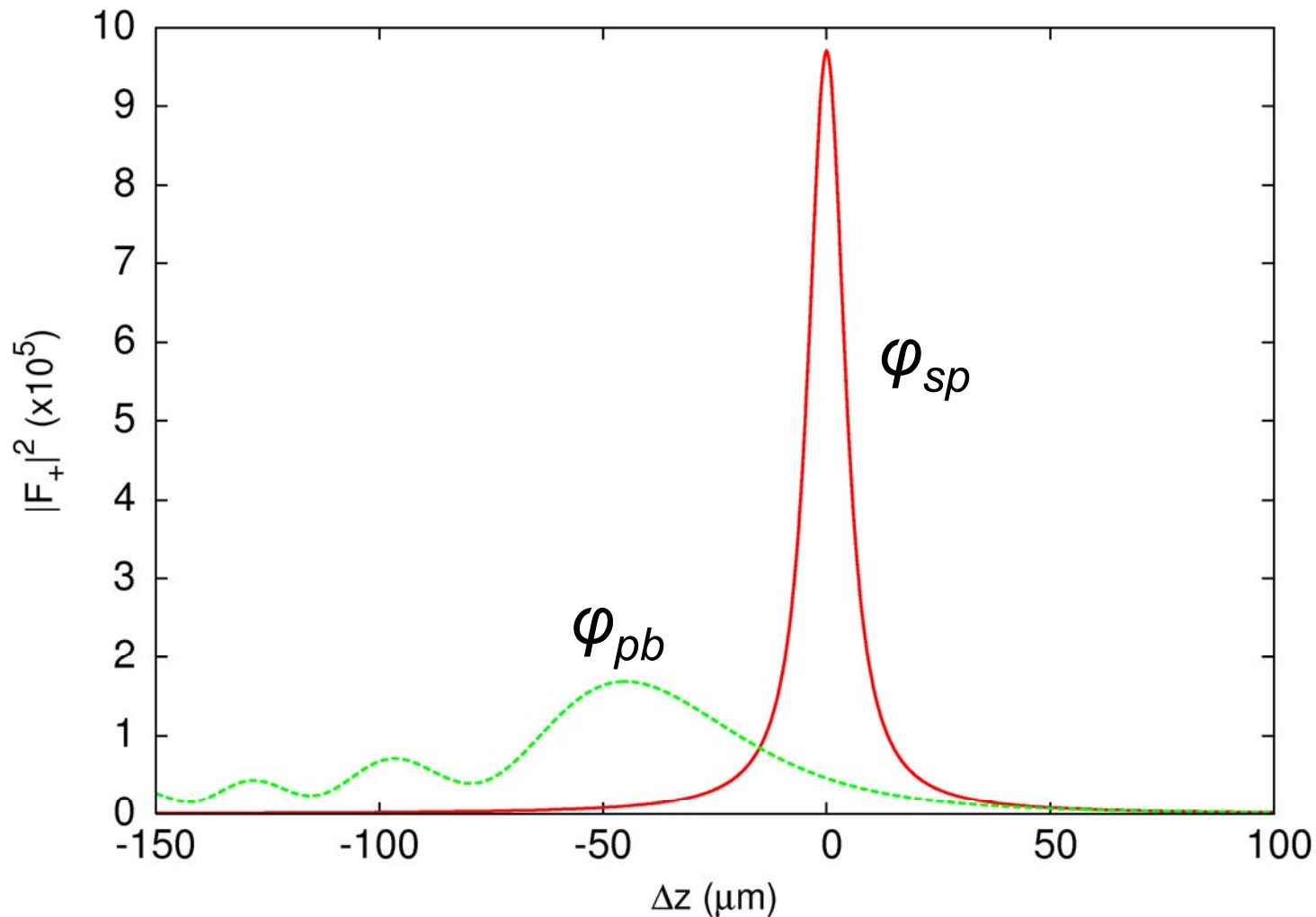
Our contribution

spherical wave front after the lens



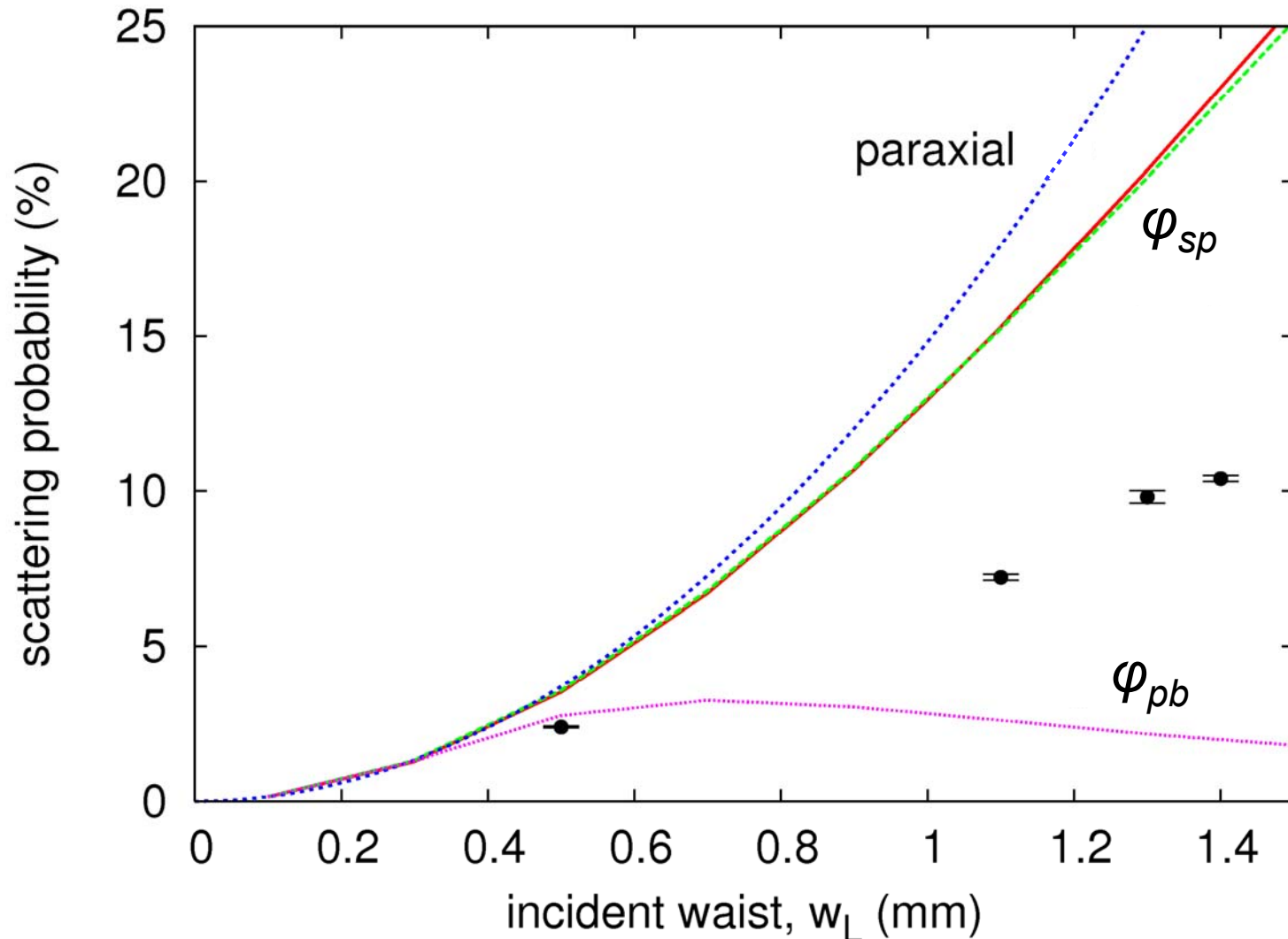
Theoretical model

Comparison of field strengths at the focus, calculated with parabolic (φ_{pb}) and spherical (φ_{sp}) phase factors for our experimental parameters



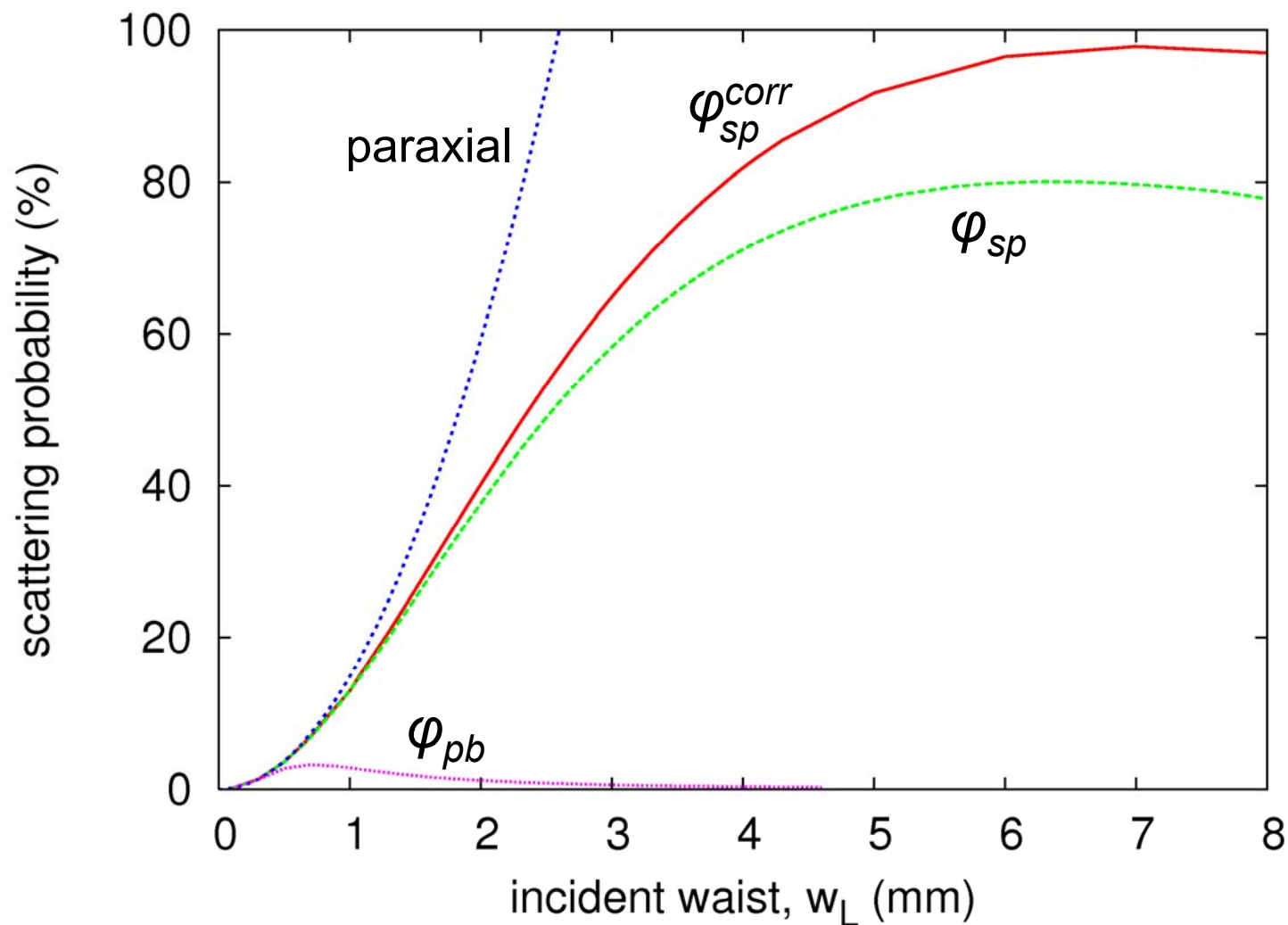
Comparison of experiment and theory

Scattering probability dependence on the input waist of the Gaussian beam for the lens used in our experiment



Theoretical model

If one focuses stronger....



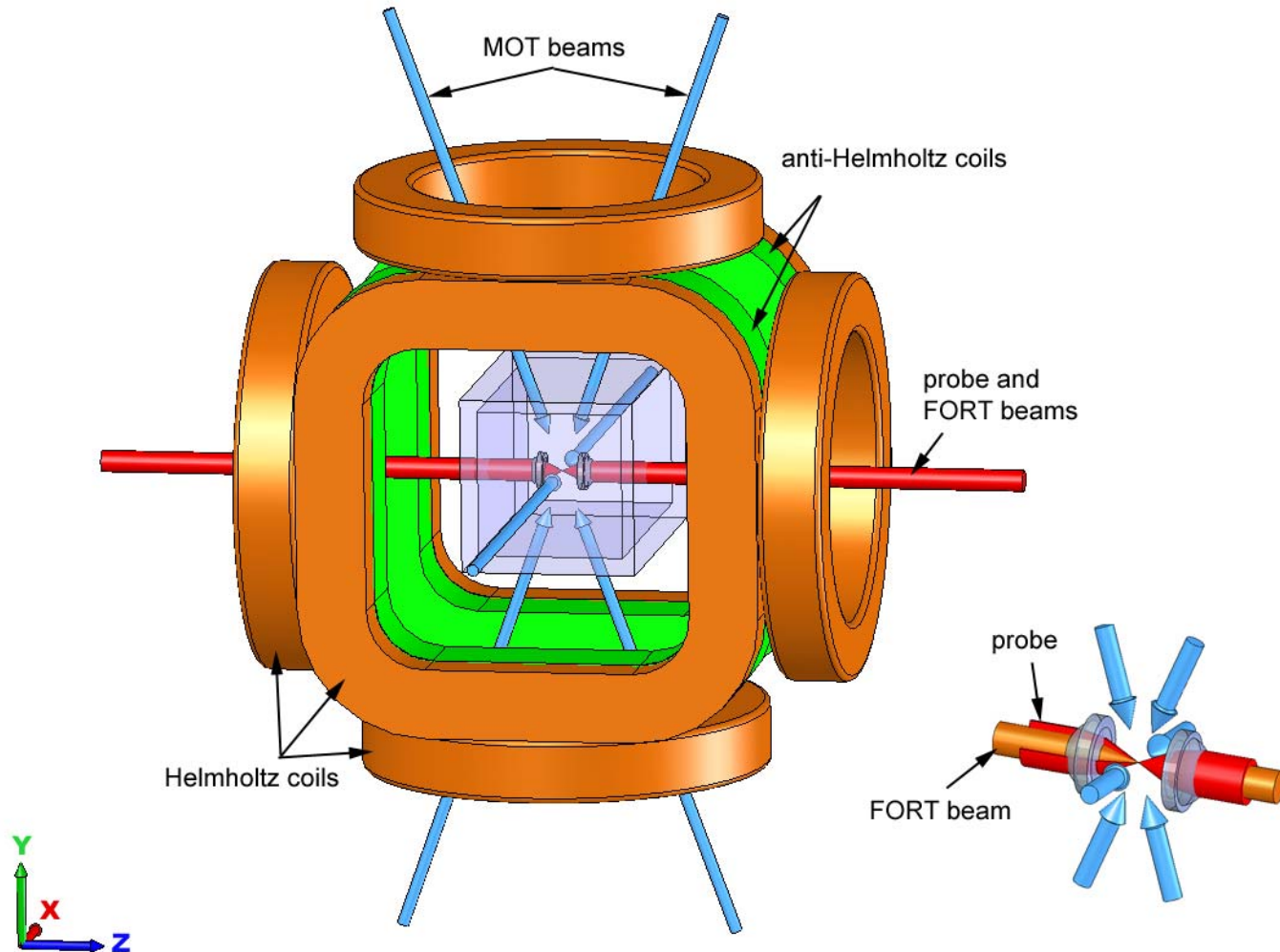
Conclusions

Substantial scattering probability can be achieved for a focused coherent light beam!

Experiment: scattering probability of **10.4 %** was directly measured for the maximum focusing achievable with current lens setup.

Theory: scattering probability up to **98%** is predicted for lenses with realistic focal length.

Thank you for your attention!



Antibunching in single atom fluorescence

